

# High-Altitude Data Assimilation System Experiments for the Northern Summer Mesosphere Season of 2007

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## Abstract

A global numerical weather prediction system is extended to mesospheric and lower thermospheric altitudes and used to assimilate high-altitude satellite measurements of temperature, water vapor and ozone from MLS and SABER during May-July 2007. Assimilated temperatures from 100-0.001 hPa show minimal biases compared to satellite data and existing analysis fields. Saturation ratios derived diagnostically from assimilated temperature and water vapor fields at PMC altitudes and latitudes compare well with seasonal variations in PMC frequency derived from measurements from the Aeronomy of Ice in the Mesosphere (AIM) satellite. Synoptic

maps of these diagnostic saturation ratios correlate geographically with three independent transient mesospheric cloud events observed at mid-latitudes by SHIMMER on STPSat-1 and by ground observers during June 2007. Assimilated temperatures and winds reveal broadly realistic amplitudes of the quasi 5-day wave and migrating tides as a function of latitude and height. The 5-day wave and migrating diurnal tide also produce water vapor responses in the polar summer MLT. These features do not correlate linearly with corresponding temperature amplitudes and thus may have a more complex origin than the 5-day wave response in PMC brightness.

*Key words:*

data assimilation, polar mesospheric clouds

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## 1 Introduction

Global numerical weather prediction (NWP) systems consist of two main components: a forecast model that predicts future atmospheric conditions, and a data assimilation system (DAS) that provides global initial conditions for those forecasts based on available observations. The quality of these operationally assimilated fields, and hence the skill of the resulting forecasts, rely heavily on the high-density planetary-scale measurements provided from satellites.

The maturity and sophistication of NWP DASs have now led them to directly assimilate satellite radiances, since the forecast models now provide more accurate *a priori* estimates than the climatologies typically used in standard

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11 satellite retrievals. Radiance weighting functions, however, are typically verti-  
12 cally broad, with some having long “tails” that extend to high altitudes not  
13 covered by the forecast model component. This latter restriction reduces the  
14 accuracy of the radiance assimilation, and can prevent certain channels from  
15 being assimilated. Thus most operational centers are progressively increasing  
16 the vertical range of their NWP systems to span most of the stratosphere,  
17 with a few now extending into the lower mesosphere.

18 To date, however, no operational NWP system extends through the meso-  
19 sphere and lower thermosphere (MLT), due primarily to a lack of operational  
20 radiance channels in this altitude range. This is changing with the launch  
21 of the latest generation of operational sensors, such as the Special Sensor  
22 Microwave Imager/Sounder (SSMIS) (Lipton, 2002; Kerola, 2006), and the  
23 advent of fast radiative transfer codes suitable for assimilating MLT radiances  
24 operationally (Han et al., 2007). Thus research NWP systems extending into  
25 the MLT are now being actively developed (Polavarapu et al., 2005a; Hoppel  
26 et al., 2008). Since the boundary between atmosphere and space is defined  
27 arbitrarily at an MLT altitude of either 80 or 100 km, such systems represent  
28 the first steps towards a truly integrated global ground-to-space forecasting  
29 capability. Here we work with one such developmental ground-to-MLT NWP  
30 system: the Advanced Level Physics High-Altitude (ALPHA) prototype of the  
31 Navy Operational Global Atmospheric Prediction System (NOGAPS), which  
32 we describe in section 2.

33 Adding an MLT component to an NWP system presents a variety of technical  
34 challenges that are only just beginning to be grappled with. The vertically  
35 extended forecast model must include new physical processes appropriate for  
36 the MLT: those for NOGAPS-ALPHA are discussed in section 2.1. Coupling

37 the forecast model to the DAS raises further issues, such as vertical informa-  
38 tion transfer, model-data biases, appropriate dynamical balance constraints,  
39 and resolved and parameterized gravity wave dynamics (Polavarapu et al.,  
40 2005b; Sankey et al., 2007). In developing and testing new MLT components,  
41 research MLT data, such as provided by the Aeronomy of Ice in the Meso-  
42 sphere (AIM) satellite (Russell et al., 2008), are particularly valuable, either  
43 for direct assimilation into the system or for independent validation of MLT  
44 DAS output.

45 Research satellites typically measure only a targeted subset of atmospheric pa-  
46 rameters specific to their core science objectives, usually within a limited range  
47 of longitudes, latitudes, heights and local times. AIM, for example, consists of  
48 the Solar Occultation for Ice Experiment (SOFIE), which performs limb oc-  
49 culation measurements of temperature and constituents at a fixed local time  
50 (Gordley et al., 2008), and the Cloud Imaging and Particle Size (CIPS) instru-  
51 ment, which images polar mesospheric cloud (PMC) properties in the vertical  
52 (McLintock et al., 2008). By contrast, NWP systems can assimilate data from  
53 a range of satellite and suborbital instruments to provide an optimal global  
54 state estimate, filling any spatial or temporal gaps using the global model’s  
55 full-physics forecast fields constrained by DAS-based initial conditions. The  
56 synoptic gridded analysis products that result are more amenable to research  
57 studies, while providing a range of other atmospheric parameters that may  
58 not be measured directly (e.g., winds). The scientific value of such DAS prod-  
59 ucts for NASA’s middle atmosphere satellite research missions is already well  
60 established. Global analysis fields generated by the Met Office DAS, for in-  
61 stance, played a pivotal role in science studies based on data from the Upper  
62 Atmosphere Research Satellite (UARS), while NASA’s Global Modeling and

63 Assimilation Office (GMAO) GEOS DAS analyses are central to all aspects of  
64 the current Aura and Aqua missions (see, e.g., Susskind et al., 2006; Manney  
65 et al., 2007). Thus one motivation for the present work is to generate synop-  
66 tic analysis fields extending into the MLT that provide analogous support for  
67 AIM.

68 There are many ways in which ground-to-MLT DAS fields could contribute  
69 to the science return from AIM. For example, there is debate as to the role  
70 of vertical and latitudinal tidal and planetary wave transport in controlling  
71 PMC variability (e.g., Berger and vonZahn, 2007; Gerding et al., 2007; Stevens  
72 et al., 2008). The limited local time coverage of satellite MLT measurements,  
73 however, presents well-known difficulties in isolating mean, tidal and fast plan-  
74 etary wave signatures. Data analysis studies to date have employed complex  
75 asynoptic mapping or least-squares fitting algorithms that require assumptions  
76 about stationarity, aliasing and seasonal dependences (e.g., Wu et al., 1995;  
77 Burrage et al., 1995; Forbes et al., 1997; Zhu et al., 2005). While these as-  
78 sumptions can be tested and the procedures improved with the help of MLT  
79 fields from general circulation models (Oberheide et al., 2003; McLandress  
80 and Zhang, 2007) and addition of data from other instruments (Drob et al.,  
81 2000; Azeem et al., 2000), final mean mesospheric temperature estimates from  
82 these algorithms can still have large uncertainties (Drob et al., 2000; Ober-  
83 heide et al., 2003; Zhu et al., 2005). NWP systems combine aspects of all the  
84 aforementioned algorithms by optimally assimilating MLT data from a variety  
85 of sources with the aid of a full-physics general circulation model (GCM) to  
86 constrain the system dynamically and optimally fill gaps. The physical and dy-  
87 namical constraints of the system yield additional benefits, such as estimates  
88 of atmospheric parameters not directly measured, such as winds. Furthermore,

89 mean and root-mean-square (RMS) differences between the model forecasts  
90 (F), observations (O) and the analysis fields (A) provide objective quantifi-  
91 cation of the inherent biases and uncertainties of all the analyzed physical  
92 quantities output by the system.

93 We explore these potential benefits for MLT science in this paper. After de-  
94 scribing the system in section 2 and tuning it in section 3, we validate its  
95 output against independent observations and analysis fields in section 4. In  
96 section 5 we study the seasonal variation of temperature and water vapor at  
97 PMC altitudes from the NOGAPS-ALPHA analysis and compare with cor-  
98 responding PMC data acquired from AIM. Section 6 studies planetary wave  
99 signals in the analysis fields, focusing on the quasi 5-day wave and solar mi-  
100 grating tides near PMC regions. Section 7 applies the synoptic analysis fields  
101 to mid-latitude mesospheric cloud events reported at specific geographical lo-  
102 cations on specific dates. Section 8 summarizes the major findings of these  
103 assimilation experiments and assesses near-term development needs for the  
104 system to improve MLT products in future assimilation experiments.

## 105 **2 NOGAPS-ALPHA**

106 Here we briefly describe the salient aspects of the NOGAPS-ALPHA system  
107 used in this study. Hoppel et al. (2008) provide a complete overview of the  
108 initial system developed for data assimilation research.

110 Hogan and Rosmond (1991) and Hogan et al. (1991) provide detailed descrip-  
111 tions of the NOGAPS global forecast model. Briefly, the dynamical core is  
112 Eulerian, hydrostatic, spectral in the horizontal and finite difference in the  
113 vertical, using the specific Lorenz-grid vertical discretization of Arakawa and  
114 Suarez (1983) generalized to hybrid vertical coordinates following Simmons  
115 and Burridge (1981). The model is forwarded using a three time-level scheme  
116 incorporating a semi-implicit treatment of gravity wave propagation, implicit  
117 zonal advection of moisture and constituents, and Robert (Asselin) time filter-  
118 ing. The operational model’s physical parameterizations include vertical diffu-  
119 sive transport in the planetary boundary layer (Louis, 1979; Louis et al., 1982)  
120 coupled to a land surface model (Hogan, 2007), orographic gravity-wave and  
121 flow-blocking drag (Webster et al., 2003), shallow cumulus mixing (Tiedtke,  
122 1984), deep cumulus convection (Peng et al., 2004), convective, stratiform and  
123 boundary layer clouds and precipitation (Slingo, 1987; Teixeira and Hogan,  
124 2002), and shortwave and longwave radiation (Harshvardhan et al., 1987).  
125 At the Fleet Numerical Meteorological and Oceanographic Center (FNMOC)  
126 NOGAPS runs operationally at T239L30 (T119L30 for ensemble forecasts)  
127 using mean orography, pure  $\sigma$  levels, a rigid upper boundary at  $p_{top} = 1$  hPa  
128 and layer thicknesses that yield a highest undiffused model layer at  $\sim 25$  hPa.

129 The progressive extension of this forecast model through the stratosphere  
130 and into the lower mesosphere for NOGAPS-ALPHA has been described by  
131 Eckermann et al. (2004) and Allen et al. (2006). We briefly summarize salient  
132 additions here, focusing mostly on new physical parameterizations required to  
133 forecast the MLT and support the DAS at these altitudes.

134 *2.1.1 Radiation*

135 The high-altitude forecast model has been designed to switch easily between  
136 any of the operational or research physics packages. Thus, while the opera-  
137 tional Harshvardhan et al. (1987) radiation schemes can be used, NOGAPS-  
138 ALPHA runs here use different schemes that extend to MLT altitudes. Radia-  
139 tive heating rates are computed using the Chou and Suarez (1999) scheme.  
140 We deactivate their near-infrared (IR) CO<sub>2</sub> band contributions at upper levels  
141 and use instead upper-level rates from Fomichev et al. (2004) that better pa-  
142 rameterize non-local thermodynamic equilibrium (non-LTE) effects on these  
143 near-IR bands. Longwave cooling rates are computed using the parameteri-  
144 zation of Chou et al. (2001), which is accurate from the ground to 0.01 hPa,  
145 and of Fomichev et al. (1998), which includes non-LTE effects on IR CO<sub>2</sub>  
146 emissions at MLT altitudes. The two profiles  $Q_{Chou}(Z)$  and  $Q_{Fomichev}(Z)$  are  
147 blended into a final cooling rate profile

148 
$$Q(Z) = w(Z)Q_{Chou}(Z) + [1 - w(Z)]Q_{Fomichev}(Z), \quad (1)$$

149 using a pressure-height dependent linear weight

150 
$$w(Z) = \frac{1 - \tanh\left(\frac{Z - Z_{int}}{\zeta}\right)}{2}, \quad (2)$$

151 where  $Z_{int} = 75$  km and  $\zeta = 5$  km. Due to the computational expense, here we  
152 update radiative heating and cooling rates every 2 hours.

153 *2.1.2 Trace Constituents*

154 Specific humidity  $q$  is built into the discretized NOGAPS primitive equations  
155 through the virtual potential temperature. Thus NOGAPS-ALPHA must ini-

156 tialize and forecast it accurately from the surface to the MLT. In addition to  
157 NOGAPS tropospheric moist physics parameterizations, we have developed  
158 new parameterizations of water vapor production in the stratosphere due to  
159 methane oxidation and photolytic loss in the mesosphere (McCormack et al.,  
160 2008). However, as the rates are generally slow except at the very highest  
161 altitudes, they were not used in the data-assimilation runs reported here, so  
162 that middle atmospheric water vapor was simply advected passively by the  
163 forecast model. Only prognostic  $q$  values below 200 hPa altitude are used in  
164 the radiation calculations: above that, values from observational and model-  
165 based climatologies are used (see section 3.1.1.1 of Eckermann et al., 2007, for  
166 details).

167 The forecast model incorporates a new prognostic capability for ozone, with  
168 a number of ozone photochemistry parameterizations available for use (Eck-  
169 ermann et al., 2004; McCormack et al., 2004). Here we use the scheme of  
170 McCormack et al. (2006) that is run operationally in the National Centers  
171 for Environmental Prediction Global Forecast System (NCEP GFS). It uses  
172 lookup tables of diurnally averaged photochemical coefficients derived from a  
173 full chemistry model based on linearizing scaled odd-oxygen production and  
174 loss rates about equilibrium states. These equilibrium states are specified in  
175 the model using zonal-mean observational climatologies, and are chosen care-  
176 fully here to match characteristics of the assimilated ozone observations so as  
177 to avoid model-data bias (Geer et al., 2007; Coy et al., 2007). The scheme does  
178 not at present parameterize either diurnal ozone photochemistry at altitudes  
179 above  $\sim 0.3$  hPa or tropospheric ozone chemistry, relaxing ozone in both alti-  
180 tude regions to a reference state based on a mean photochemical relaxation  
181 rate. Because of this, here we do not use prognostic ozone mixing ratios  $\chi_{O_3}$  in

182 the radiation calculations. Instead we use the observational ozone climatology  
 183 described by Eckermann et al. (2007) that incorporates daytime ozone data  
 184 only at high altitudes. That climatology was improved slightly here by adding  
 185 high-altitude daytime ozone data from the High Resolution Doppler Imager  
 186 (Marsh et al., 2002) at the very highest altitudes. These daytime ozone clima-  
 187 tologies improve the model’s radiative heating rates (Eckermann et al., 2007)  
 188 which are important for accurately modeling and assimilating temperature at  
 189 these altitudes (e.g., Sassi et al., 2005). For simplicity, the night-to-day ra-  
 190 tio used by Eckermann et al. (2007) to scale up nighttime ozone values for  
 191 the cooling rate calculations was not used since it is a more minor effect for  
 192 temperature prediction.

### 193 *2.1.3 Gravity Wave Drag*

194 Nonorographic gravity wave drag (GWD) is the most important new param-  
 195 eterization required for summer MLT prediction (e.g., Fritts and Luo, 1995).  
 196 Here we use a multiwave scheme based on the linear GW saturation formula-  
 197 tion of Lindzen (1981), as developed for the Whole Atmosphere Community  
 198 Climate Model (WACCM). Appendix A of Garcia et al. (2007) provides a de-  
 199 tailed description. Here we summarize the important aspects for the present  
 200 work.

201 The scheme launches a prescribed spectrum of  $n_{gw}$  individual GWs within  
 202 every grid box at a source level set here to 500 hPa (following Garcia et al.,  
 203 2007). Each GW  $j$  is assigned a unique ground-based horizontal phase speed

$$\begin{aligned}
 c_j &= |U_{500}| + j\Delta c, \\
 j &= -n_c, -n_c + 1, \dots, +n_c - 1, +n_c, (j \in \mathbf{Z}, n_c \in \mathbf{N}),
 \end{aligned}
 \tag{3}$$

204 such that phase speeds are distributed symmetrically with respect to the  
 205 500 hPa horizontal wind speed  $|U_{500}|$ . As in Garcia et al. (2007) we choose  
 206  $\Delta c = 2.5 \text{ m s}^{-1}$  and  $n_c = 32$ , yielding  $n_{gw} = 2n_c + 1 = 65$  component gravity  
 207 waves with intrinsic phase speeds  $|c_j - U_{500}|$  distributed between  $\pm 80 \text{ m s}^{-1}$ ,  
 208 all aligned along the 500 hPa wind speed direction. The vertical flux of hor-  
 209 izontal pseudomomentum density (Eliassen-Palm flux) of each wave,  $\tau_{src}(c_j)$ ,  
 210 is assigned based on a Gaussian flux distribution versus phase speed, centered  
 211 about  $|U_{500}|$ , of the form

$$212 \quad \tau_{src}(c_j) = \tau_b F(\phi, t) \exp \left[ \frac{-(c_j - c_{j=0})^2}{\hat{c}_w^2} \right]. \quad (4)$$

213 Following Garcia et al. (2007) we set the phase-speed width  $\hat{c}_w = 30 \text{ m s}^{-1}$  and  
 214 tune the so-called background flux  $\tau_b$  in experiments described in section 3.  
 215 The function  $F(\phi, t)$ , plotted in Figure 1, is an analytical fit as a function  
 216 of latitude  $\phi$  and time (month)  $t$  to results obtained from diagnostically pro-  
 217 cessing long-term climate model output using a proposed parameterization of  
 218 frontogenetic gravity wave generation (c.f. Figure 2 of Charron and Manzini,  
 219 2002). It yields a large winter-summer flux asymmetry in each hemisphere  
 220 as well as a gradual variation during the season: see Garcia et al. (2007) for  
 221 further details.

222 Wave propagation, wave breaking and saturation, and resulting diffusive and  
 223 thermal dissipation of wave momentum flux are modeled for each wave as in  
 224 Garcia et al. (2007). At model layer  $k$ , the ensuing GW-induced mean-flow  
 225 acceleration due to wave  $j$  is

$$226 \quad a_{j,k} = -g\epsilon \frac{\partial \tau_{j,k}}{\partial p}, \quad (5)$$

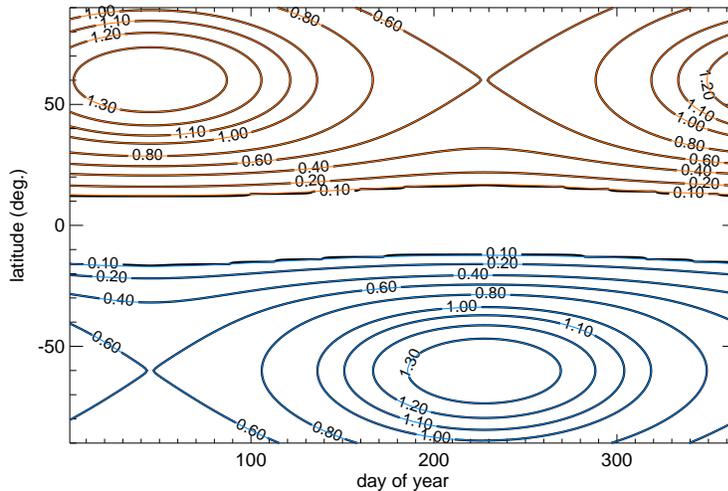


Fig. 1. GWD source function  $F(\phi, t)$ . Orange and blue contours show its northern and southern hemisphere components, respectively.

227 where  $g$  is gravitational acceleration and  $p$  is pressure. This acceleration, di-  
 228 rected along the 500 hPa (source-level) wind speed direction, is apportioned  
 229 to modify zonal and meridional wind speeds accordingly, and summed over  
 230 all GWs  $j$ . All remaining flux is deposited in the top two model layers to con-  
 231 serve momentum so as to capture robust downward-control responses (Shaw  
 232 and Shepherd, 2007).

233 The factor  $\epsilon$  in (5) is a constant set in the range  $0 \leq \epsilon \leq 1$ . Such normaliza-  
 234 tion terms occur routinely in Lindzen-type schemes to ameliorate excessively  
 235 large and/or insufficiently smooth GWD in models, and are usually justified  
 236 physically as encapsulating either the net “efficiency” of wave breaking or the  
 237 net “intermittency” of GW activity, either in time (due to variable forcing)  
 238 or spatially due to incomplete filling of the grid box by the GW packets. The  
 239 mathematical implementation of the efficiency concept in (5), which follows  
 240 that used in the Alexander and Dunkerton (1999) scheme, simply scales down  
 241 all the GWD values by a constant amount. It should be noted that similar

242 efficiency factors are implemented in different ways mathematically in other  
243 Lindzen schemes, where they act differently and modify both the shape and  
244 magnitude of the GWD profile (Hamilton, 1997; McLandress, 1998; Norton  
245 and Thuburn, 1999). The specific implementation here has the practical ad-  
246 vantage that source flux parameters like  $\tau_b$  can be adjusted to modify the  
247 shape of the GWD profile, whereupon  $\epsilon$  can then be adjusted to scale the final  
248 GWD while retaining the tuned profile shape. Specific GWD tuning for the  
249 NOGAPS-ALPHA assimilation runs is discussed in section 3.

250 For the experiments reported here, we apply only the scheme’s GW momentum  
251 flux divergence tendencies to the model: GWD-induced vertical diffusivities,  
252 while calculated, are not at present used to mix momentum, heat or con-  
253 stituents in the model. Orographic GWD is applied separately using either a  
254 Palmer et al. (1986) or Webster et al. (2003) scheme: we choose the former  
255 here following Siskind et al. (2007).

#### 256 *2.1.4 Resolution and Height Range*

257 As in Hoppel et al. (2008) the forecast model is run here at a triangular  
258 spectral truncation of T79, corresponding to  $1.5^\circ$  longitude resolution on the  
259 quadratic Gaussian grid. We use 68 model layers that extend into the MLT  
260 ( $p_{top} = 5 \cdot 10^{-4}$  hPa) with a vertical pressure height resolution  $\Delta Z \approx 2$  km  
261 throughout the middle atmosphere. The runs here use the “NEWHYB2” hy-  
262 brid vertical coordinate described by Eckermann (2008) with  $k_p = 43$  isobaric  
263 model layers between  $p_{top}$  and  $p_{k_p+1/2} \sim 87.4$  hPa. This new hybrid coordinate  
264 reduces vertical truncation errors in the stratosphere and MLT (Eckermann,  
265 2008) and should improve the quality of the assimilations above the tropopause

266 (see, e.g., Trenberth and Stepaniak, 2002).

## 267 2.2 DAS Component

268 The NOGAPS-ALPHA DAS uses a three-dimensional variational (3DVAR)  
269 algorithm formulated in observational space, known as the Naval Research  
270 Laboratory (NRL) Atmospheric Variational DAS, or NAVDAS. The basic for-  
271 mulation and initial performance are described by Daley and Barker (2001a),  
272 while Daley and Barker (2001b) provide a more detailed description. Hoppel  
273 et al. (2008) explain how NAVDAS was interfaced to the NOGAPS-ALPHA  
274 forecast model to run as a coupled NWP system extending into the MLT. The  
275 version and setup used here are very similar to those described by Hoppel et  
276 al. (2008), and so we focus here mainly on an overview of that system and  
277 salient differences for the experiments reported here.

278 Given a column vector  $\mathbf{x}_b$  containing  $I$  “background” estimates of some atmo-  
279 spheric parameter (e.g., temperature) and  $L$  estimates of some related obser-  
280 vation  $\mathbf{y}$  (e.g., thermal radiance), NAVDAS generates a corresponding analysis  
281 vector  $\mathbf{x}_a$  by numerically minimizing the scalar cost function

$$282 \quad J(\mathbf{x}_a) = (\mathbf{y} - \mathcal{H}(\mathbf{x}_a))^T \mathbf{R}^{-1} (\mathbf{y} - \mathcal{H}(\mathbf{x}_a)) + (\mathbf{x}_b - \mathbf{x}_a)^T \mathbf{P}_b^{-1} (\mathbf{x}_b - \mathbf{x}_a). \quad (6)$$

283 Here  $\mathcal{H}$  is the forward observation operator (e.g., a radiative transfer forward  
284 model that converts temperature into radiance) and  $\mathbf{R}$  is the  $L \times L$  error  
285 covariance matrix of this conversion: similarly  $\mathbf{P}_b$  is the  $I \times I$  error covariance  
286 matrix in the background estimate, and  $T$  denotes transpose. (6) is identical  
287 in form to cost functions solved in standard satellite retrievals except here the  
288 background  $\mathbf{x}_b$  is provided by the forecast model.

289 The observation-space solution to (6) is (Daley and Barker, 2001a)

$$290 \quad \mathbf{x}_a - \mathbf{x}_b = \mathbf{P}_b \mathbf{H}^T [\mathbf{H} \mathbf{P}_b \mathbf{H}^T + \mathbf{R}] [\mathbf{y} - \mathcal{H}(\mathbf{x}_b)], \quad (7)$$

291 which converts the so-called innovations  $\mathbf{y} - \mathcal{H}(\mathbf{x}_b)$  in the observation space  
292 into a correction vector  $\mathbf{x}_a - \mathbf{x}_b$  in model/analysis space. The matrix  $\mathbf{H} =$   
293  $\partial \mathcal{H} / \partial \mathbf{x} |_{\mathbf{x}_b}$  originates as an approximation of  $\mathcal{H}(\mathbf{x}_a)$  by the truncated Taylor  
294 series expansion  $\mathcal{H}(\mathbf{x}_b) + \mathbf{H}[\mathbf{x}_a - \mathbf{x}_b]$ . The accuracy of this approximation, and  
295 hence the quality of the analysis, clearly requires *inter alia* that  $\mathbf{x}_a - \mathbf{x}_b$  be as  
296 small as possible, and thus that the forecast model be minimally biased with  
297 respect to both observations and analysis (see section 3).

298 Research forecast-assimilation runs with NOGAPS-ALPHA use the archived  
299 sensor data routinely assimilated operationally by NOGAPS at FNMOC (see  
300 Table 1 of Baker et al., 2007). Most relevant here for the middle atmosphere  
301 are Advanced Microwave Sounding Unit (AMSU-A) thermal radiances from  
302 stratospheric channels 9 and 10 on the NOAA-15 and NOAA-16 satellites.  
303 We use the NAVDAS operational data-thinning schemes for these radiances  
304 described by Baker et al. (2005).

305 In the initial NOGAPS-ALPHA implementation, Hoppel et al. (2008) assim-  
306 ilated limb-scanned temperature data from:

- 307 • the Microwave Limb Sounder (MLS) on Aura from 32-0.01 hPa (version 2.2  
308 retrievals: see Schwartz et al., 2008);
- 309 • the Sounding of the Atmosphere Using Broadband Emission Radiometry  
310 (SABER) instrument on the Thermosphere Ionosphere Mesosphere Ener-  
311 getics and Dynamics (TIMED) satellite from 32-0.019 hPa (version 1.06  
312 retrievals: see Mertens et al., 2004).

313 In this study we also assimilate SABER and MLS temperatures using the same  
314 observation operators and error covariances for these instruments described by  
315 Hoppel et al. (2008). Here, however, we use version 1.07 SABER retrievals that  
316 account for the vibrational exchange between CO<sub>2</sub> isotopes. Kutepov et al.  
317 (2006) have shown that this process is critical for getting reliable temperatures  
318 in the summer mesopause.

319 Here we assimilate SABER and MLS temperatures up to a higher altitude of  
320 0.002 hPa, in order to insert data at polar summer mesopause altitudes. As  
321 in Hoppel et al. (2008) the increments above this top data insertion level are  
322 progressively damped over a ~6 km pressure height range, before reverting  
323 thereafter to pure forecast fields. Hoppel et al. (2008) fitted and removed  
324 a global mean profile of the bias between SABER and MLS temperatures,  
325 using the latter as truth to bias correct the SABER data. The calculation is  
326 repeated here for the May-June period using the version 1.07 SABER data  
327 over this higher altitude range. The resulting bias profile in Figure 2 is similar  
328 to independent estimates of both Hoppel et al. (2008) and Schwartz et al.  
329 (2008).

330 PMC formation and microphysics depend sensitively on not just temperature  
331 but water vapor abundances, which in turn depend on HO<sub>x</sub>/O<sub>x</sub> chemistry.  
332 Thus in these experiments we also assimilate version 2.2 MLS retrievals of  
333  $\chi_{H_2O}$  and  $\chi_{O_3}$ .

334 Lambert et al. (2007) and Read et al. (2007) provide detailed descriptions  
335 and validation of the version 2.2 MLS  $\chi_{H_2O}$  retrievals at upper and lower  
336 altitudes, respectively. They recommend that science studies using these data  
337 be confined to the 316-0.002 hPa range. Here we assimilate MLS  $\chi_{H_2O}$  profiles

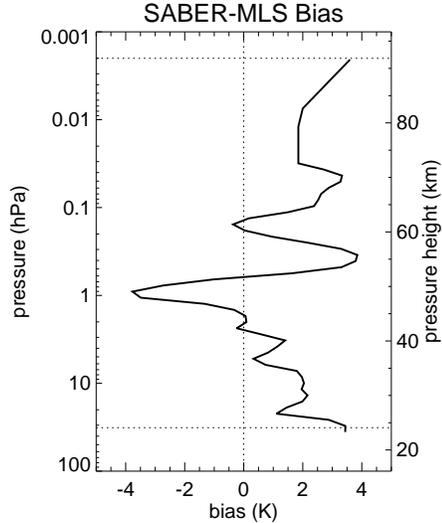


Fig. 2. Global-mean bias of SABER using MLS as truth, estimated using O-F statistics from the test assimilation (15 May-20 June) following Hoppel et al. (2008). Upper and lower limits of data insertion are marked with dotted lines.

338 from 220-0.002 hPa. We use horizontal correlation lengths of 358 km for these  
 339 data, consistent with consensus values in Table 2 of Lambert et al. (2007).  
 340 The effective vertical resolution of these data increase with altitude but, as  
 341 for MLS temperatures, we use a constant mean vertical Gaussian averaging  
 342 with a full-width half-maximum (FWHM) of  $\sim 4$  km: more detailed discussion  
 343 of these choices is provided in section 3.2 of Hoppel et al. (2008).

344 We assimilate MLS  $\chi_{O_3}$  profiles from 215-0.02 hPa, the same validated range  
 345 quoted in Table 1 of Jiang et al. (2007), who showed good agreement with sub-  
 346 orbital ozone profiles throughout the middle atmosphere. Since the forecast  
 347 model does not parameterize diurnal ozone photochemistry (see section 2.1),  
 348 we assimilate only daytime values at altitudes above 1 hPa. The vertical reso-  
 349 lution of the data is  $\sim 3$  km throughout the range, which we use as our vertical  
 350 averaging FWHM for these data in the DAS along with the same horizontal  
 351 correlation lengths as for MLS temperature of  $\sim 380$  km.

352 The system runs in a standard 6-hour forecast-assimilation cycle, which presents  
353 challenges for assimilating tides (Swinbank et al., 1999): given their impor-  
354 tance in the MLT, the following choices were made to aid assimilation of tidal  
355 features. The standard nonlinear normal mode initialization (NNMI) of the  
356 analysis state by the forecast model was deactivated, given that it mishandles  
357 migrating tides (Wergen, 1989). NOGAPS-ALPHA can perform NNMI on  
358 the analysis increments only, which treats tides at lower altitudes much better  
359 while still eliminating gravity wave noise (Ballish et al., 1992; Seaman et al.,  
360 1995). However, since its potential impact at new MLT altitudes has not been  
361 methodically investigated, we opted to deactivate it too and thus perform no  
362 initial-state filtering of the analysis prior to running forecasts. Work by Sankey  
363 et al. (2007) with the Canadian Middle Atmosphere Model (CMAM) suggests  
364 that the additional resulting GW noise in the forecast can propagate into the  
365 MLT and break, affecting mean and tidal structures, with some damping of  
366 tidal amplitude and spreading of tidal frequencies: nonetheless, their work in-  
367 dicates that broadly realistic diurnal and semidiurnal tides are still captured  
368 in the MLT analysis fields.

### 369 **3 Reduction of Model Bias By Tuning the GWD Parameterization**

370 Biases in model forecasts yield biased analysis fields (Dee and daSilva, 1998).  
371 Thus, prior to commencing the assimilation runs reported here, we performed  
372 an iterative series of 2-week forecasts that were initialized to high-altitudes  
373 at various times in June 2007 using output from a one-month test assimila-  
374 tion. The forecasts of zonal-mean temperature were compared with time series  
375 of MLS and SABER zonal-mean temperatures at various latitudes and pres-

376 sures. These comparisons led to adjustment of the background flux  $\tau_b$  and/or  
377 efficiency  $\epsilon$  used in the nonorographic GWD scheme (see section 2.1.3), where-  
378 upon the forecasts and comparisons were repeated until the forecast temper-  
379 atures were close to the observations throughout June, focusing especially on  
380 the polar summer MLT.

381 The top four panels in Figure 3 summarize results of four different forecast  
382 experiments each initialized on 1 June, showing zonal-mean summer hemi-  
383 sphere temperatures after +14 days. Forecasts in Figure 3a used the default  
384  $\tau_b$  and  $\epsilon$  settings of Hoppel et al. (2008) and yield a polar summer mesopause  
385 that is located too low in altitude, followed by an unrealistically sharp tem-  
386 perature gradient yielding a thin warm layer at  $\sim 0.002$  hPa: zonal-mean MLS  
387 temperatures on 15 June 2007 are shown for reference in Figure 3f. Reduc-  
388 ing  $\tau_b$  by a factor of 4 to 1.75 mPa yields forecast temperatures in Figure 3b  
389 that all but eliminate the secondary warm layer and generate a polar sum-  
390 mer mesopause at roughly the right altitude, but which is too warm relative  
391 to MLS and SABER. Increasing  $\epsilon$  by a factor of 2 yields an excessively cold  
392 mesopause (Figure 3c). An intermediate choice of  $\epsilon = 0.0175$  yields zonal-mean  
393 forecast temperatures in Figure 3d with a polar summer mesopause of about  
394 the right altitude and temperature according to both SABER and MLS. Fore-  
395 casts initialized at other times in June produce similarly good results using  
396 these settings (see, e.g., Figure 3e).

397 Changes in  $\tau_b$  and  $\epsilon$  above apply globally to reduce GWD in the winter (south-  
398 ern) hemisphere as well, where they yield forecast temperatures that compare  
399 less favorably with MLS and SABER due to reduced diabatic descent. Thus we  
400 performed additional forecasts that retained the tuned summer GWD settings  
401 in Figure 3d but increased winter GWD by scaling up  $F(\phi, t)$  in the south

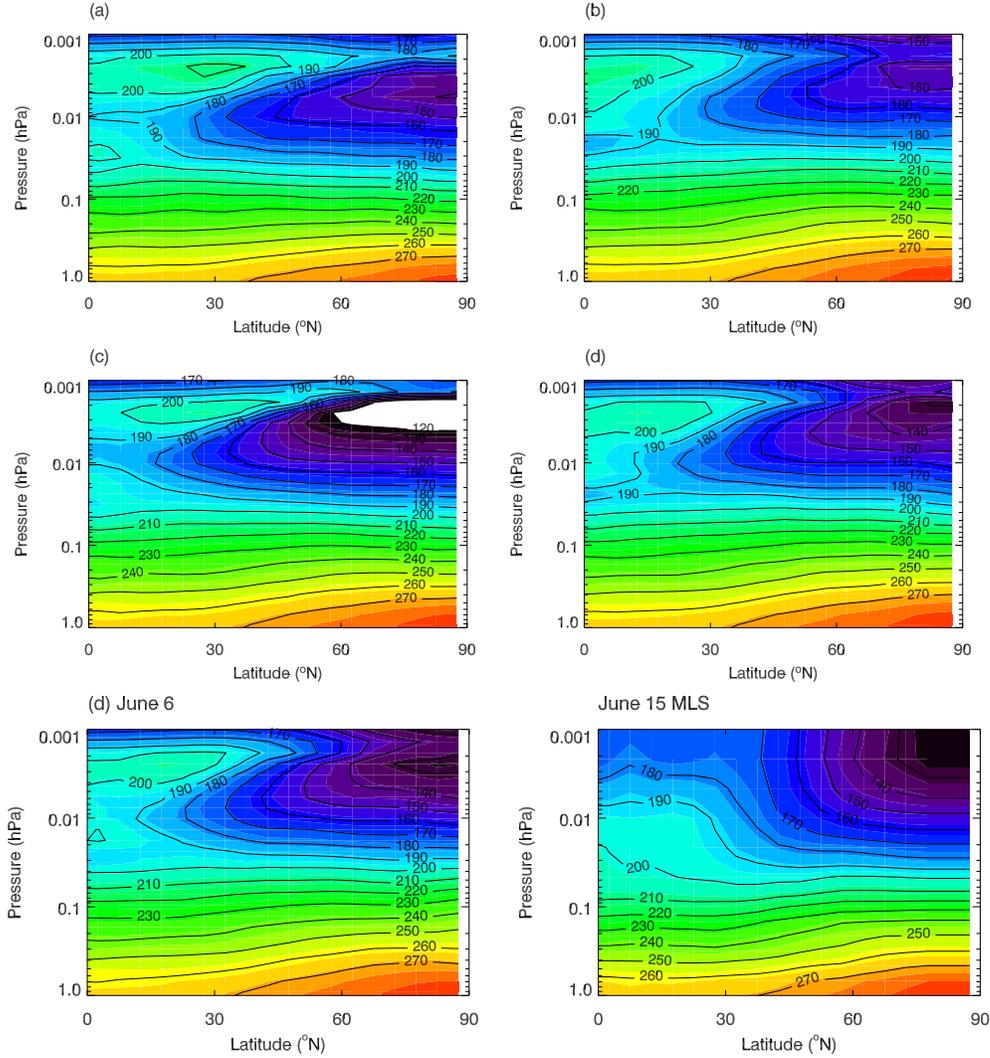


Fig. 3. Zonal-mean northern-hemisphere temperatures from 1-0.001 hPa of four NOGAPS-ALPHA +14 day forecasts, initialized on 1 June 2007 using preliminary high-altitude analysis fields and the following nonorographic GWD parameter settings: (a)  $\tau_b = 7$  mPa,  $\epsilon = 0.0125$ ; (b)  $\tau_b = 1.75$  mPa,  $\epsilon = 0.0125$ ; (c)  $\tau_b = 1.75$  mPa,  $\epsilon = 0.0250$ ; (d)  $\tau_b = 1.75$  mPa,  $\epsilon = 0.0175$ . Panel (e) shows results using same settings as (d) but initialized on 6 June. Panel (f) shows zonal-mean MLS temperatures on 15 June. Temperatures below 120 K are not plotted.

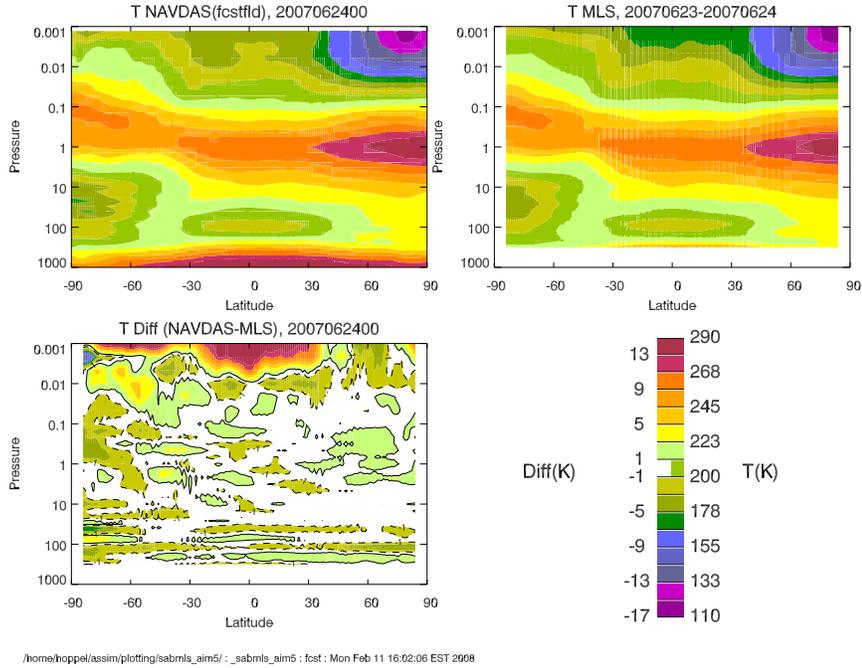
402 (blue curve in Figure 1). A series of these experiments (not shown) yielded  
 403 summer MLT forecasts consistently poorer than Figure 3d, due to the indi-  
 404 rect effects of increased winter GWD on summer MLT temperatures through

405 a modified mesospheric residual circulation (Becker and Fritts, 2006). Thus we  
406 settled here upon tuned GWD settings of  $\tau_b = 1.75$  mPa and  $\epsilon = 0.0175$  for our  
407 final assimilation run to reduce the potential for large mean O-F (innovation)  
408 biases in the summer MLT, with the understanding that forecast biases may  
409 be more significant in the winter hemisphere in these runs.

#### 410 **4 Initial Validation**

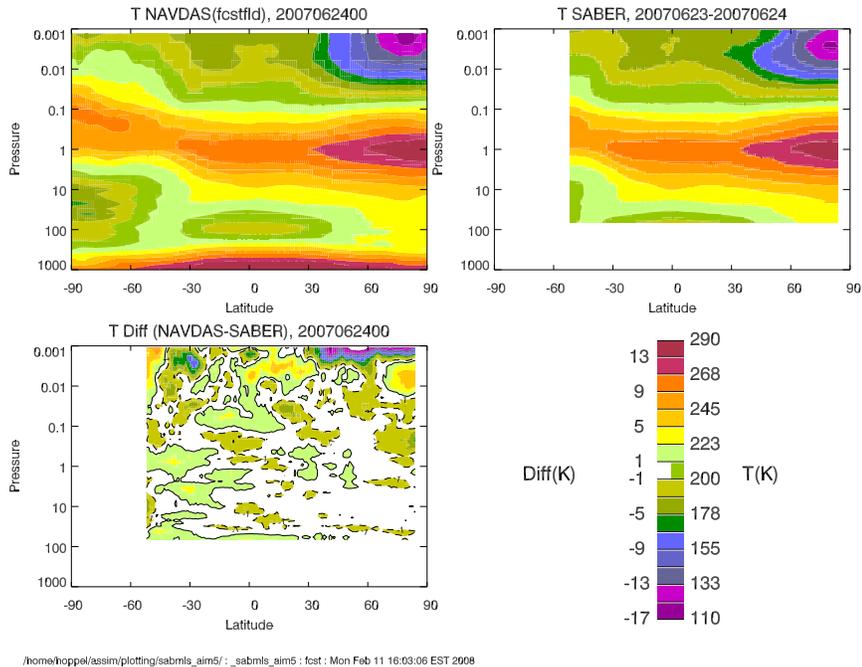
411 The final assimilation run that was subsequently performed extends over  
412 nearly a full PMC season, from 15 May 2007 to 10 August 2007. However,  
413 on 15 July the TIMED satellite yawed so that SABER no longer measured  
414 the polar summer MLT. Thus, hereafter we analyze results only up to  $\sim 15$   
415 July, during which both MLS and SABER were each contributing data to  
416 the polar summer MLT assimilation. The system generates regularly gridded  
417 global analysis fields of geopotential heights, temperatures, water vapor and  
418 ozone mixing ratios, horizontal winds, and other quantities every 6 hours at  
419 60 reference pressure levels distributed roughly evenly in pressure height over  
420 the range 1000–0.0005 hPa.

421 As the assimilation proceeded, quality checks were performed by comparing  
422 zonal-mean temperature output with separate MLS and SABER zonal means  
423 computed from 2 days of data. While MLS and SABER each provide good  
424 global coverage over 2 days, their local time sampling is limited and different.  
425 Figures 4 and 5 plot examples of such comparisons for 24 June 2007. Differ-  
426 ences are plotted in the lower panel of these figures, and are all generally small  
427 at altitudes below  $\sim 0.01$  hPa.



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Fig. 4. Zonal-mean temperatures from (a) NOGAPS-ALPHA analysis on 24 June 2007, (b) MLS on 23-24 June 2007, and (c) differences between (a) and (b).



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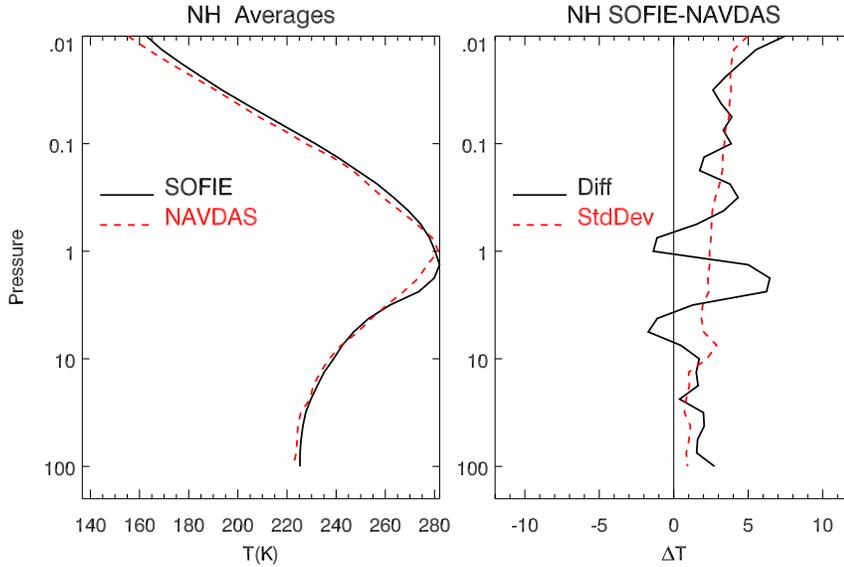
Fig. 5. Same presentation as Figure 4 but using SABER temperatures instead of MLS.

428 At higher altitudes Figure 4 shows a warm bias relative to MLS at the equator  
 429 and subtropics that is not seen in the SABER comparison in Figure 5. The

430 feature is too broad latitudinally and too stationary in time to be explained by  
431 aliasing of the diurnal tide, and seems to reflect systematic biases between the  
432 two measurements, highlighting potential weaknesses in our use of a single bias  
433 correction profile in Figure 2. Since MLS averaging kernels become vertically  
434 broad here, our use of constant vertical averaging widths for MLS data in  
435 the DAS may also be problematic. Additionally, the forecast model does not  
436 reproduce a realistic semiannual oscillation of the equatorial MLT with the  
437 current GWD settings due to insufficient tropical GW flux, an issue currently  
438 being addressed separately by tropical GWD tuning experiments over climate  
439 time scales with the forecast model.

440 The SABER comparisons in Figure 5 show a systematic warm bias at the  
441 uppermost polar summer MLT altitudes that is not seen in the MLS compar-  
442 isons. Kutepov et al. (2006) note that a residual warm bias in version 1.07  
443 SABER temperatures may still exist above  $\sim 86$  km altitude in polar summer.  
444 However, the bias here could instead reflect cold biases in MLS at these alti-  
445 tudes like those seen at other latitudes in Figure 4. If borne out, then future  
446 bias correction schemes may be needed that apply separately to both MLS  
447 and SABER. Overall, however, these comparisons reveal consistency among  
448 MLS, SABER and analyzed temperatures at most altitudes, including at and  
449 just below summer mesopause altitudes where PMCs form.

450 Next we compared assimilated fields with preliminary retrievals from SOFIE.  
451 These comparisons serve two purposes. First, as measurements not assimilated  
452 into the system, SOFIE data provide independent validation of the analy-  
453 sis. Second, SOFIE retrievals are at an early stage of development, and can  
454 themselves benefit from validation studies. Thus, these comparisons should be  
455 viewed as mutual cross-validation of two independent emerging AIM-related



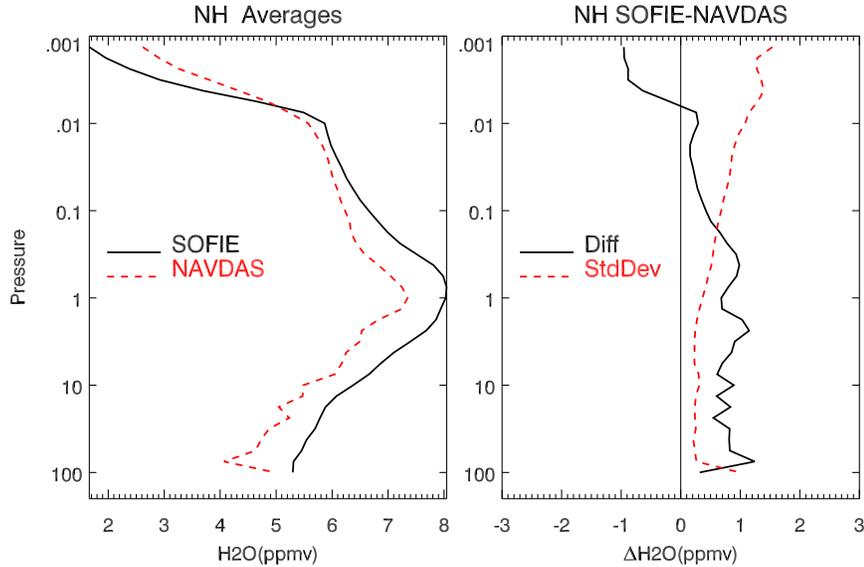
Dates: 20070525 - 20070623, Number of profiles = 433, Lat range: ~66N - 69N

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Fig. 6. (a) Mean temperatures from 25 May-23 June 2007 between  $66^{\circ}$ – $69^{\circ}$ N from 433 retrieved SOFIE profiles (black) and the NOGAPS-ALPHA analysis at the closest longitude, latitude and time to the SOFIE measurement (red); (b) corresponding mean bias (black) and standard deviation (red) between the two.

456 products.

457 Figure 6a compares the mean SOFIE temperatures from the northern hemi-  
 458 sphere between  $66^{\circ}$ – $69^{\circ}$ N from 25 May to 23 June with analyzed temperature  
 459 profiles from NOGAPS-ALPHA at the nearest longitude, latitude and time of  
 460 each measurement (433 profiles in all). The comparisons reveal that SOFIE  
 461 temperatures are already very close to the MLS and bias-corrected SABER  
 462 temperatures assimilated by NOGAPS-ALPHA from 100-0.01 hPa. Figure 6b  
 463 plots the corresponding mean bias and standard deviation, which are both gen-  
 464 erally small and increase slowly with altitude. The bias increases with height  
 465 may be due to small errors in the SOFIE retrieval, while standard deviation  
 466 increases with height may be related to small-scale noise in the NOGAPS-



Dates: 20070525 - 20070623, Number of profiles = 433, Lat range: ~66N - 69N

/home/hoppel/assim/plotting/sabmls\_aim5/ : \_sabmls\_aim5 : Wed Feb 6 12:21:29 EST 2008 : /home/hoppel/assim/sofie/sofie\_nh.sav

Fig. 7. Same presentation as in Figure 6 but for  $\chi_{H_2O}$  (in ppmv).

467 ALPHA forecasts.

468 Figure 7 shows the corresponding plot for  $\chi_{H_2O}$ , which reveals a systematic  
 469 SOFIE wet bias throughout the stratosphere and lower mesosphere with re-  
 470 spect to assimilated MLS water from NOGAPS-ALPHA. The high bias in  
 471 these early experimental SOFIE  $\chi_{H_2O}$  retrievals results from channel align-  
 472 ment biases caused by pressure registration and field-of-view offsets which  
 473 are in the process of being corrected (Gordley, private communication, 2008).  
 474 This comparison highlights the usefulness of NOGAPS-ALPHA as an early  
 475 validation standard for emerging retrieval products from AIM instruments.

476 Finally we compare zonal-mean zonal winds, temperatures and ozone mix-  
 477 ing ratios with lower-altitude DAS products from a more mature system:  
 478 the NASA GEOS4 (Bloom et al., 2005). While the GEOS4 runs extend to  
 479 0.01 hPa, the standard analysis fields are issued only to 0.2 hPa as shown in  
 480 Figure 8b. The overall June 2007 morphology of the mean zonal wind jets and

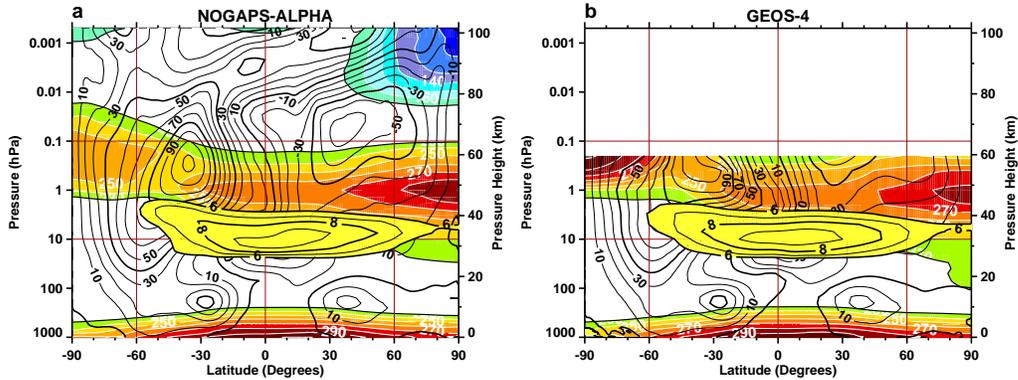


Fig. 8. Comparison of zonal-mean analysis output for June 2007 from (a) NOGAPS-ALPHA and (b) NASA GEOS4, showing zonal winds ( $\text{m s}^{-1}$ , black contours), temperatures (K, rainbow color scale and white contours), and ozone mixing ratios (ppmv, black contours in foreground yellow).

481 stratopause temperature structure is very similar in both analyses, apart from  
 482 temperature differences near the polar winter stratopause. While the GWD  
 483 parameter settings in NOGAPS-ALPHA may yield forecast biases in the win-  
 484 ter hemisphere (see section 3), the NOGAPS-ALPHA polar winter stratopause  
 485 temperatures compare better to MLS than the GEOS4 values (note that the  
 486 TIMED satellite's yaw cycle had SABER preferentially viewing high northern  
 487 latitudes at this time: see Figure 5). The zonal-mean peak ozone mixing ratios  
 488 near 10 hPa are also very similar between the two analyses. The NOGAPS-  
 489 ALPHA results at higher altitudes in Figure 8a show closure of the extratrop-  
 490 ical mesospheric zonal wind jets in both hemispheres and a cold mean polar  
 491 summer mesopause. We now look in more detail at the polar summer MLT as  
 492 described by the NOGAPS-ALPHA analysis fields.

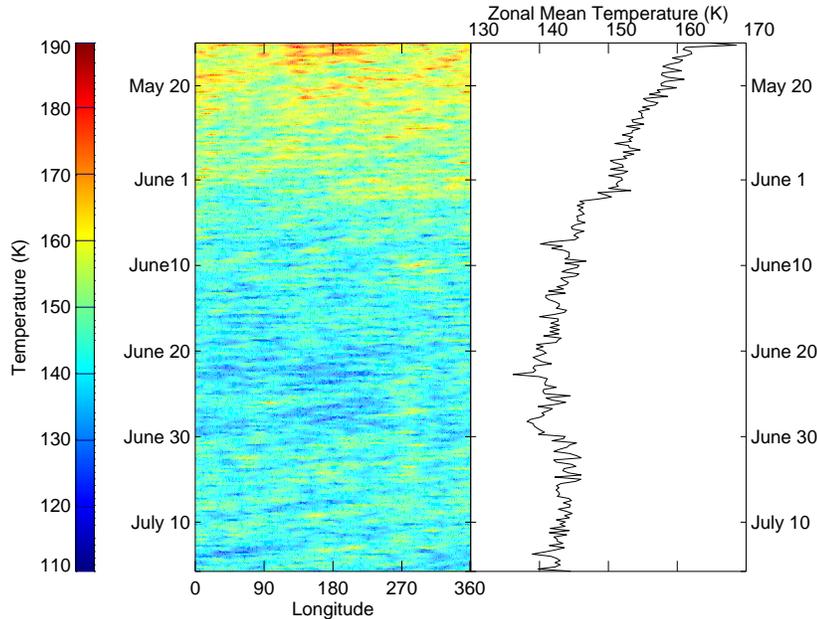


Fig. 9. Plot of mean temperatures between  $65^{\circ}$ – $70^{\circ}$ N at 0.006 hPa from 15 May-15 July 2007, shown in Hovmöller form on the left and the corresponding zonal-mean temperature time series on the right.

## 493 5 Mean Variations During PMC Season

494 Next we look at the mean polar summer MLT thermal conditions relevant to  
 495 PMCs as provided from the analysis fields.

496 Figure 9 plots the time variation of temperature at 0.006 hPa (a typical PMC  
 497 altitude) averaged between  $65^{\circ}$ – $70^{\circ}$ N. The analysis captures a gradual march  
 498 to lower temperatures through May and June, yielding cold values in late June  
 499 and mid July. The Hovmöller and zonal-mean plots both show spatial variability  
 500 and temporal intermittency on both large and small scales throughout the  
 501 season. Some of this can immediately be seen to be geophysical. For example,  
 502 in late July a wavenumber-1 oscillation with a period of  $\sim 5$  days is evident.  
 503 Quasi 5-day waves in temperatures and other analyzed parameters are studied

504 in greater depth in section 6.

505 Next we combine NOGAPS-ALPHA temperatures  $T$  and water vapor mixing  
506 ratios  $\chi_{H_2O}$  to derive diagnostic saturation ratios for ice,

$$507 \quad S = \frac{p_{H_2O}}{p_{ice}}, \quad (8)$$

508 at PMC altitudes, as follows. At a given analysis pressure  $p$ , the partial pres-  
509 sure of water vapor  $p_{H_2O} = p \chi_{H_2O}$ . We specify the saturation vapor pressure  
510 for ice,  $p_{ice}$ , using the Murphy and Koop (2005) fit

$$511 \quad \log p_{ice} = 9.550426 - \frac{5723.265}{T} + 3.53068 \log T - 0.00728332 T, \quad (9)$$

512 which is valid for  $T \geq 110$  K, and thus valid for PMC studies (Rapp and  
513 Thomas, 2006).

514 We have computed these diagnostic  $S$  values at PMC altitudes and compared  
515 them with the various indicators of cloud occurrence frequency derived from  
516 SOFIE data that are discussed by Stevens et al. (2008). One such comparison  
517 is summarized in Figure 10, which plots time series of zonal-mean diurnally-  
518 averaged NOGAPS-ALPHA saturation ratios  $S$  at 0.006 hPa in the 65°-70°N  
519 band. The dashed curve shows corresponding time series in this latitude band  
520 of PMC occurrence frequency derived from SOFIE data, taken from Figure 5  
521 of Stevens et al. (2008). Except for the two large saturation spikes in the  
522 NOGAPS-ALPHA record, the overall agreement between the two time series  
523 is excellent, with an overall linear correlation coefficient of 0.78. Particularly  
524 noteworthy is the decrease in SOFIE PMC frequency around June 30, which  
525 is also seen in CIPS data (Merkel et al., this issue). This coincides with an  
526 abrupt  $\sim 5$  K increase in the zonal-mean temperature at 0.006 hPa at these

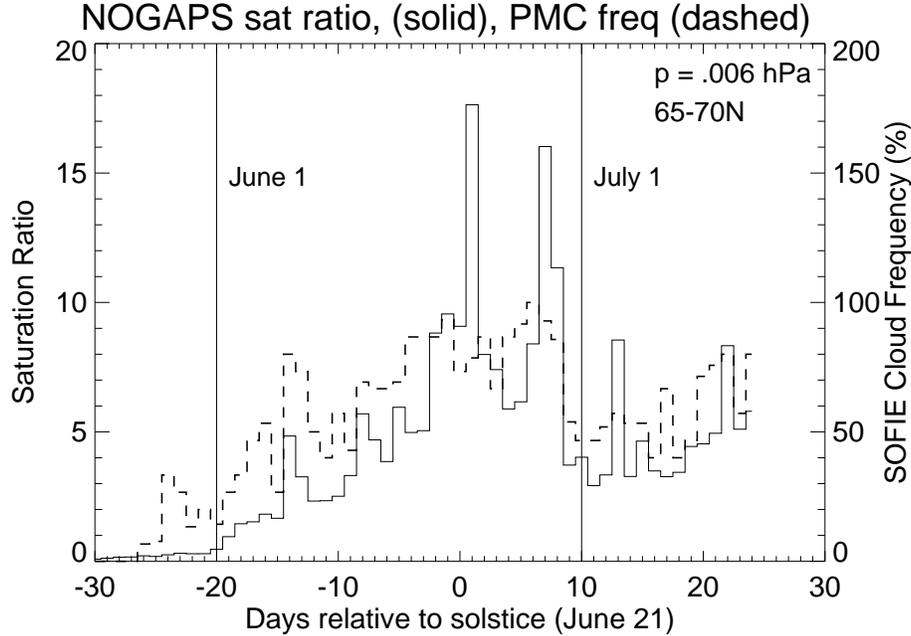


Fig. 10. NOGAPS-ALPHA saturation ratio  $S$  (solid curve, left hand axis) at 0.006 hPa between  $65^{\circ}$ – $70^{\circ}$ N compared with PMC occurrence frequency of bright clouds observed by SOFIE (dashed curve, right hand axis) taken from Figure 5 of Stevens et al. (2008). The linear correlation coefficient between the two curves is 0.78.

527 latitudes at the end of June, as seen in Figure 9.

## 528 6 Planetary Waves

529 As discussed in the introduction, accurate extraction of tides and fast plan-  
 530 etary waves from asynoptic satellite data alone is difficult. Analysis systems  
 531 like NOGAPS-ALPHA offer a potentially powerful tool for improving data-  
 532 based planetary wave estimates, given that the forecast model ideally forecasts  
 533 these waves and thus the DAS can optimally blend wave signals in both the  
 534 observations  $y$  and a forecast background  $x_b$  constrained by analyzed initial  
 535 conditions (Swinbank et al., 1999). Here we demonstrate the capability by per-

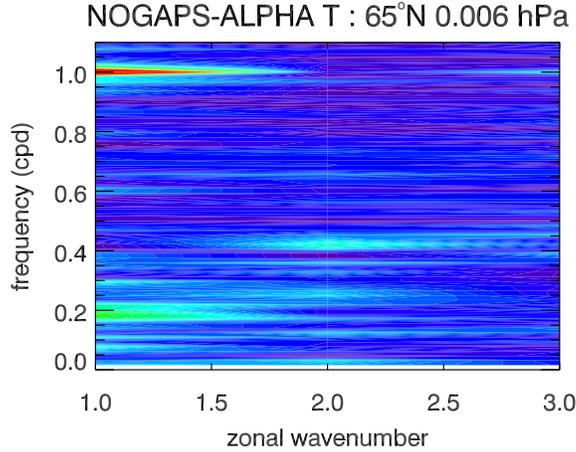


Fig. 11. Mean space-time magnitude of amplitude spectrum of temperature at 0.006 hPa and 65°N for westward propagating wavenumbers 1–3 over period 15 May to 18 July, 2007. Significant peaks (light blue, green and red) occur at wavenumber 1 at  $\sim 5$  days and 1 day, and at wavenumber 2 at  $\sim 2$  days. Color scale is linear.

536 forming space-time spectral analysis of the 6-hourly NOGAPS-ALPHA fields  
 537 (Hayashi, 1982) at a range of altitudes to infer global amplitudes of particular  
 538 planetary wave motions relevant to the polar summer MLT and PMCs (for  
 539 further details on the algorithms used, see McCormack et al., 2008).

540 Figure 11 plots the mean space-time temperature spectrum of westward-  
 541 propagating disturbances at 0.006 hPa and 65°N from 15 May 2007 to 18  
 542 July 2007. It shows peaks at wavenumber 1 at  $\sim 5$  days due to the westward-  
 543 propagating (1,1) Rossby normal mode and at 1 day due to the migrating solar  
 544 diurnal tide. We analyze these strong planetary wave signals in greater depth  
 545 in what follows. Analysis of the  $\sim 2$  day wavenumber-2 Rossby normal mode  
 546 (Merkel et al., 2008), which also appears (albeit more weakly) in Figure 11, is  
 547 left for future studies.

## 548 6.1 Quasi 5-Day Wave

549 A number of studies have reported modulations in PMC occurrence by the  
550 quasi 5-day (1,1) Rossby normal mode (e.g., Kirkwood et al., 2002; Merkel et  
551 al., 2003; von Savigny et al., 2007; Merkel et al., 2008). Figure 12 shows mean  
552 values of peak temperature and meridional wind amplitude for June 2007,  
553 derived by scaling by  $\sqrt{2}$  the RMS spectral power at westward-propagating  
554 wavenumber 1 over the 4.4-6.2 day period band. Inferred temperature ampli-  
555 tudes peak at midlatitudes and are weak at the equator, a latitudinal struc-  
556 ture in broad agreement with previous modeling and observations of this mode  
557 (Hirota and Hirooka, 1984; Riggin et al., 2006). The peak amplitudes are gen-  
558 erally weak with monthly-mean values in Figure 12a of  $\sim 1\text{--}3$  K in the summer  
559 MLT, in the range of previous observational estimates (e.g., von Savigny et  
560 al., 2007). Meridional wind responses, an indirect product of the assimilation,  
561 peak in Figure 12b primarily at the poles.

## 562 6.2 Solar Migrating Tides

563 Figure 13 plots peak amplitudes of the migrating diurnal tide averaged over the  
564 entire June 2007 analysis period, derived from spectral signals at westward-  
565 propagating wave-1 in a narrow frequency band centered at  $1 \text{ day}^{-1}$ . Fig-  
566 ure 13a reveals an equatorial temperature peak of  $\sim 5\text{--}10$  K maximizing at  
567 lower altitudes slightly south of the equator, and a meridional wind peak of  
568  $\sim 25 \text{ m s}^{-1}$  at  $\sim 25^\circ\text{S}$  latitude. Both results are in fairly good agreement with  
569 long-term data-validated CMAM results for June (Figure 2 of McLandress,  
570 2002). Temperature amplitudes also agree well with amplitudes inferred from

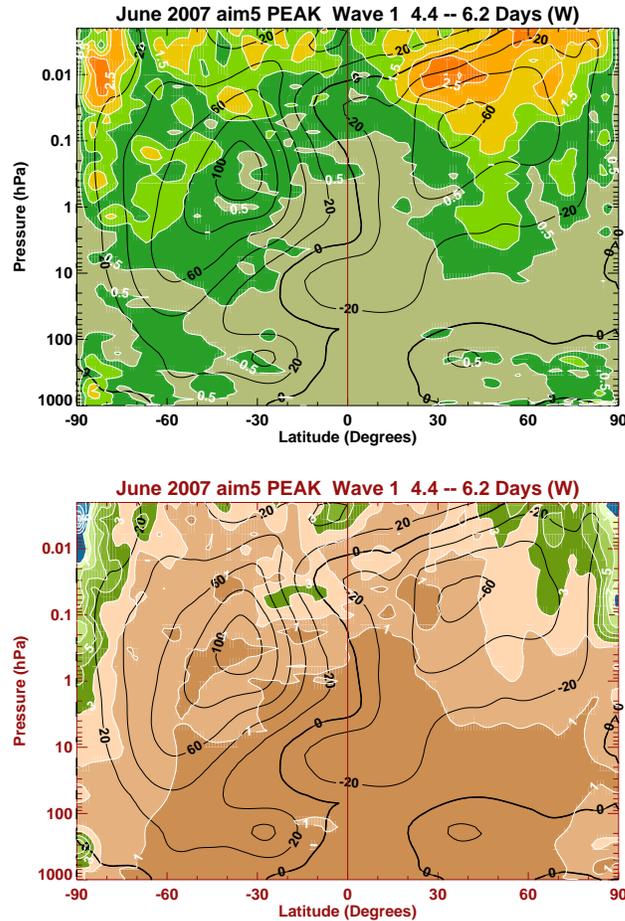


Fig. 12. White contours with associated color shading show peak amplitudes of the quasi 5-day wave computed over the 4.4-6.2 day period band from NOGAPS-ALPHA analyses for all of June 2007 in (a) temperature (Kelvin) and (b) meridional wind ( $\text{m s}^{-1}$ ). Black contours on each panel show zonal-mean zonal winds ( $\text{m s}^{-1}$ ) averaged over same period.

571 SABER temperatures during June 2004 by Zhang et al. (2006) and from UARS  
 572 MLS temperatures by Forbes et al. (2006). At higher altitudes tidal peaks of  
 573  $\sim 5$  K also occur at extratropical summer MLT altitudes.

574 Figure 14 shows corresponding results for the migrating semidiurnal tide. Ac-  
 575 curate assimilation of semidiurnal tides is particularly challenging, since they  
 576 lie at the Nyquist period of our 6 hourly 3DVAR analysis window. Moreover,

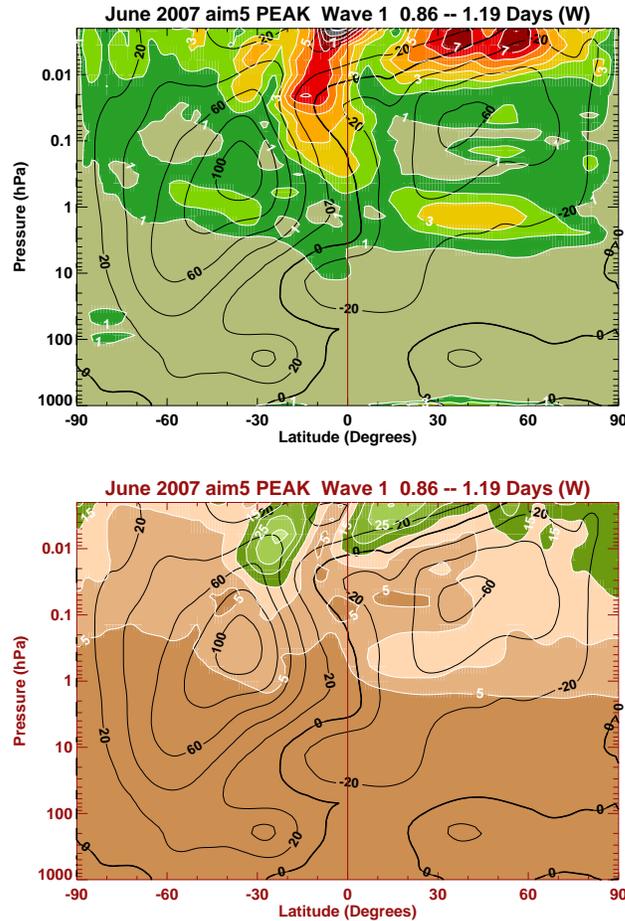


Fig. 13. White contours with associated color shading show peak diurnal tidal amplitudes from NOGAPS-ALPHA for all of June 2007 in (a) temperature (Kelvin) and (b) meridional wind ( $\text{m s}^{-1}$ ). Black contours on each panel show zonal-mean zonal winds ( $\text{m s}^{-1}$ ) averaged over same period.

577 the periodic quasi-diurnal variation in the 6-hourly longitudinal window for  
 578 MLS and SABER data insertion essentially migrates with the Sun and may  
 579 act as an artificial forcing term. Thus these semidiurnal results merit careful  
 580 scrutiny.

581 Peak temperature amplitudes in Figure 14a show an extended high-altitude  
 582 temperature peak at  $\sim 40^\circ\text{S}$  which agrees broadly with the predictions of tidal  
 583 models (see Zhang et al., 2006). The June 2004 SABER results of Zhang et

584 al. (2006) show a secondary peak nearer the equator, and our analyses show  
585 a number of similar secondary peaks at higher altitudes. In the summer MLT  
586 the temperature amplitudes are weaker in broad agreement with observations  
587 (e.g. Singer et al., 2003).

588 Our semidiurnal meridional-wind amplitudes in Figure 14b show two broad  
589 extratropical peaks peaking at  $\pm 50\text{--}60^\circ$  latitude of up to  $\sim 20\text{--}30\text{ m s}^{-1}$ . Long-  
590 term radar measurements of winds in the Arctic MLT show climatological peak  
591 meridional wind amplitudes  $\sim 10\text{--}20\text{ m s}^{-1}$  (Portnyagin et al., 2004). However,  
592 MLT winds measured by a meteor radar at Kühlungsborn ( $54^\circ\text{N}$ ) in June 2007  
593 reveal stronger semidiurnal wind amplitudes of  $\sim 20\text{--}30\text{ m s}^{-1}$  (see Figure 9 of  
594 Stevens et al., 2008) that appear to validate our indirectly inferred amplitudes  
595 of up to  $25\text{ m s}^{-1}$  at this latitude in Figure 14b.

### 596 *6.3 Planetary Wave Signals in Water Vapor*

597 Figure 15 shows the quasi 5-day wave and diurnal tidal amplitude responses  
598 in NOGAPS-ALPHA assimilated water vapor fields. The 5-day wave ampli-  
599 tudes in Figure 15a show a broad peak in the summer MLT peaking at  $\sim 60^\circ\text{--}$   
600  $75^\circ\text{N}$ . The peak amplitudes are  $\sim 0.2\text{--}0.3\text{ ppmv}$  in the lower mesosphere. These  
601 findings are consistent with ground-based microwave water vapor data from  
602 ALOMAR ( $69^\circ\text{N}$ ) analyzed by Sonnemann et al. (2008), which show quasi  
603 5-day oscillations of similar magnitude at a range of mesospheric altitudes in  
604 summer. They used a global model to explain these features in terms of 5-  
605 day wave-modulated horizontal transport across mean latitudinal water vapor  
606 gradients. Similarly located water vapor signals are seen for the diurnal tide  
607 in Figure 15b which, again, may indicate water vapor changes associated with

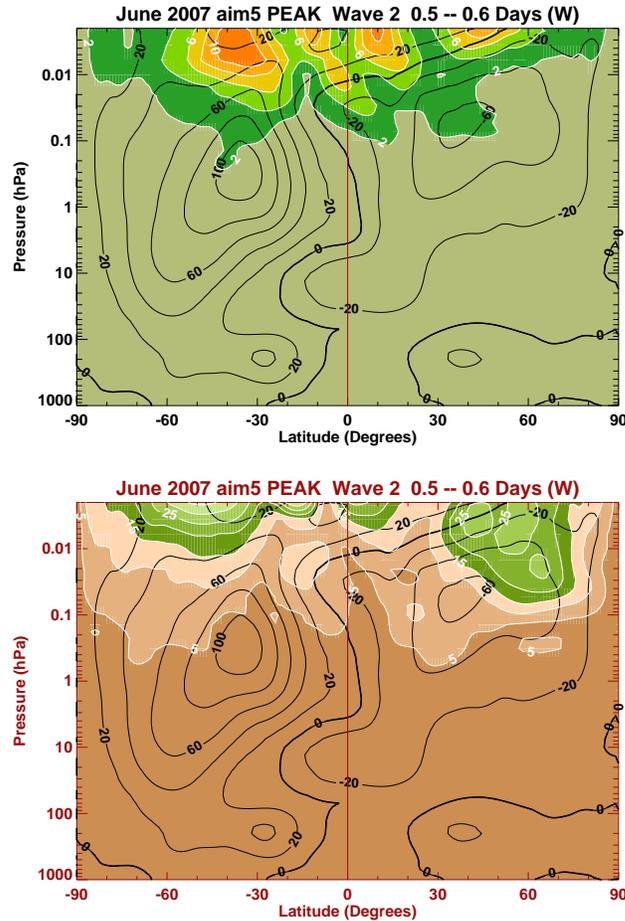


Fig. 14. As for Figure 13 but showing migrating semidiurnal tidal amplitudes.

608 diurnal tide-induced variations in meridional transport (e.g., Gerding et al.,  
 609 2007).

610 In terms of PMCs, von Savigny et al. (2007) and Merkel et al. (2008) report  
 611 strong anticorrelation between 5-day wave signals in temperature and PMC  
 612 brightness. To investigate this for water vapor, Figure 16 plots the time evo-  
 613 lution of the quasi 5-day wave amplitudes in NOGAPS-ALPHA temperature  
 614 and water vapor at the 0.006 hPa level near nominal PMC altitudes. This  
 615 was obtained by first performing a two-dimensional Fast Fourier Transform  
 616 (2DFFT) of the fields, digitally filtering the spectral components to isolate  
 617 westward zonal wavenumber 1 within the 4.4-6.25 day period band, performing

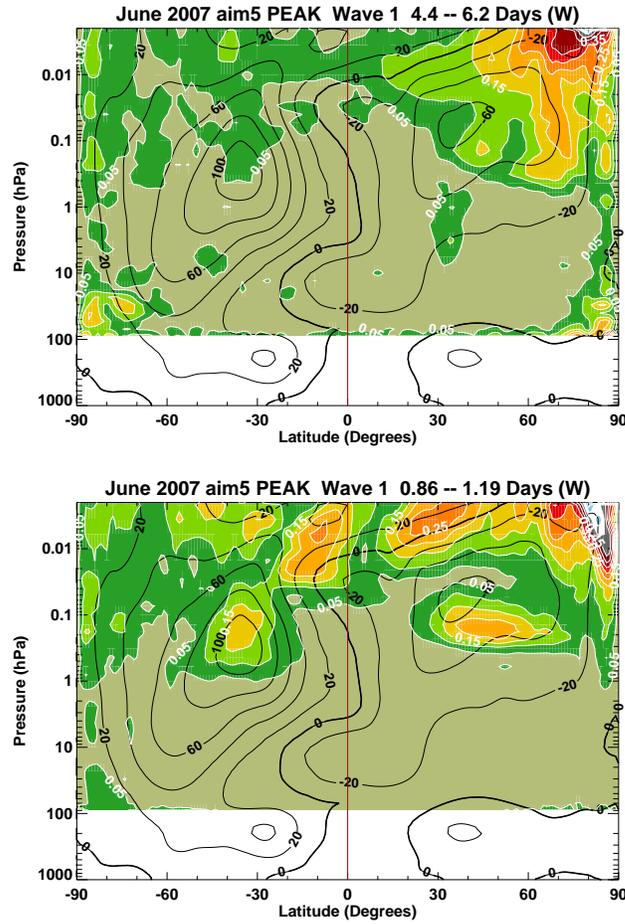


Fig. 15. White contours with associated color shading show peak amplitudes in water vapor (ppmv) for (a) quasi 5-day wave and (b) migrating diurnal tide for June 2007. Black contours on each panel show zonal-mean zonal winds ( $\text{m s}^{-1}$ ) averaged over same period.

618 an inverse 2DFFT back to the space-time domain, then computing oscillation  
 619 amplitude around a latitude circle at each latitude and time.

620 The temperature amplitudes in Figure 16a show considerable time variation,  
 621 periodically intensifying then disappearing during May and June, and reaching  
 622 largest amplitudes in early July. The corresponding water vapor amplitudes  
 623 show a weak correlation with these temperature amplitude vacillations, and  
 624 show largest amplitudes and greater day-to-day variability at polar latitudes

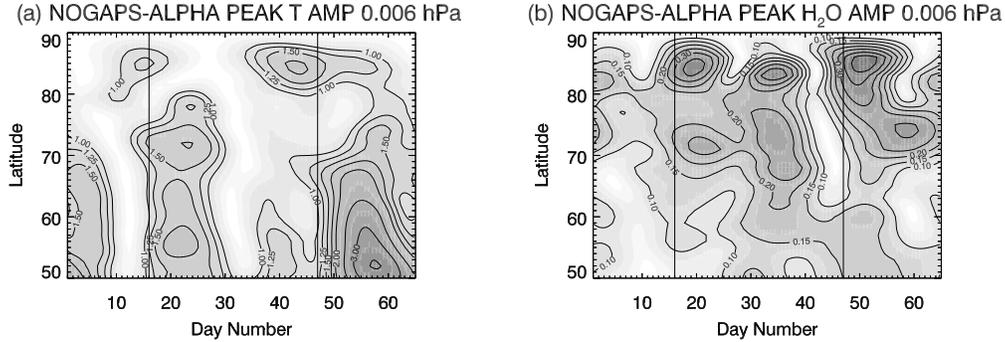


Fig. 16. Time series versus northern latitude of quasi 5-day wave peak amplitude at 0.006 hPa in (a) temperature (K), and (b) water vapor mixing ratio (ppmv). Time series starts on 15 May (day 1) and ends on 18 July (day 65). First day of June and July is marked with solid vertical line.

625 where 5-day wave meridional wind amplitudes peak (see Figure 12b). No ob-  
 626 vious correlation between 5-day wave temperature and water vapor signals is  
 627 evident in Figure 16, suggesting that the 5-day wave in water vapor mixing  
 628 ratio near PMC altitudes is not as tightly coupled as the temperature and  
 629 PMC-brightness responses at 5 days appear to be. Thus the 5-day wave signal  
 630 in water vapor mixing ratios in Figure 16b may be a more complex net effect of  
 631 temperature-modulated PMC microphysics, photolytic water vapor loss and  
 632 horizontal transport across mean  $\chi_{H_2O}$  gradients.

## 633 7 Sporadic Mid-Latitude Mesospheric Clouds

634 Equatorward of 50° latitude, mesospheric clouds (MCs) are considered by  
 635 many as a novelty, where they are more popularly referred to as noctilucent  
 636 clouds (NLCs: Taylor et al., 2002; Thomas et al., 2003). They are of global im-  
 637 portance, however, because the transition region in an environmental system is  
 638 considered to be the most sensitive to change. Theoretical studies consistently

639 indicate that climate trends should manifest themselves first in mid-latitude  
640 MCs (Thomas, 1996; Siskind et al., 2005).

641 The MC events that occur episodically at middle latitudes are not well under-  
642 stood (see, e.g., Herron et al., 2007). The global synoptic perspective afforded  
643 by NOGAPS-ALPHA assimilations is well suited to identification and analy-  
644 sis of such events. Here we illustrate the potential of NOGAPS-ALPHA fields  
645 to cast light on these mysterious events, leaving the deeper interpretation of  
646 the fundamental thermal and dynamical processes that seeded them for sub-  
647 sequent research.

#### 648 *7.1 MC event of 13 June 2007*

649 The Spatial Heterodyne Imager of Mesospheric Radicals (SHIMMER) on  
650 STPSat-1 (Space Test Program Satellite-1) is a limb-viewing UV imager. The  
651 limb view geometry makes it well suited to detecting dimmer PMCs at lower  
652 latitudes since, unlike nadir sounders, the cloud signal does not need to be  
653 discriminated from a bright background.

654 Figure 17a plots all 236 UV limb profiles from 66–100 km altitude acquired by  
655 SHIMMER on 13 June 2007 between 40°–50°N. After removing the Rayleigh  
656 background using the technique described by Stevens et al. (2008), three pro-  
657 files on this day, marked in red in Figure 17a, indicate presence of a meso-  
658 spheric cloud at  $\sim 80$  km (Figure 17b).

659 The geographical locations of these three profiles are shown with red dots in  
660 Figure 18c. Successive panels plot maps of the NOGAPS-ALPHA tempera-  
661 ture, water vapor mixing ratios and saturation ratios  $S$  at 0.006 hPa over the

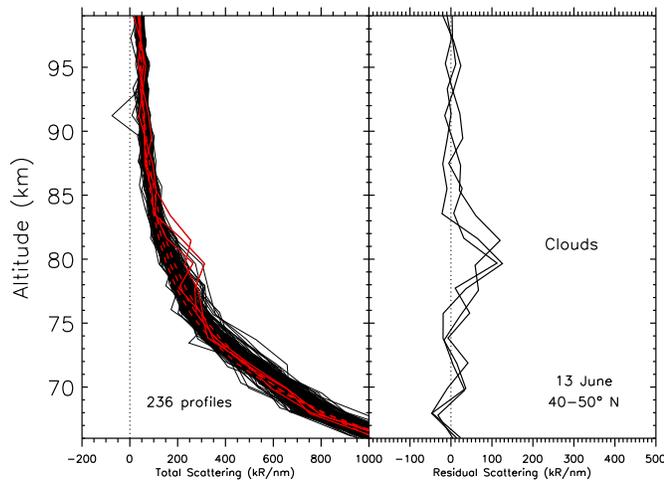


Fig. 17. SHIMMER mesospheric cloud detections on June 13, 2007. (a) All 236 limb brightness profiles taken between  $40^{\circ}$ – $50^{\circ}$ N on that day (black). The profiles are dominated by Rayleigh scattering. Scattering signal from mesospheric clouds is superimposed on this background contribution in three profiles plotted in red. (b) Difference between observed signals in (a) and fitted Rayleigh background for the three MC detections in (a).

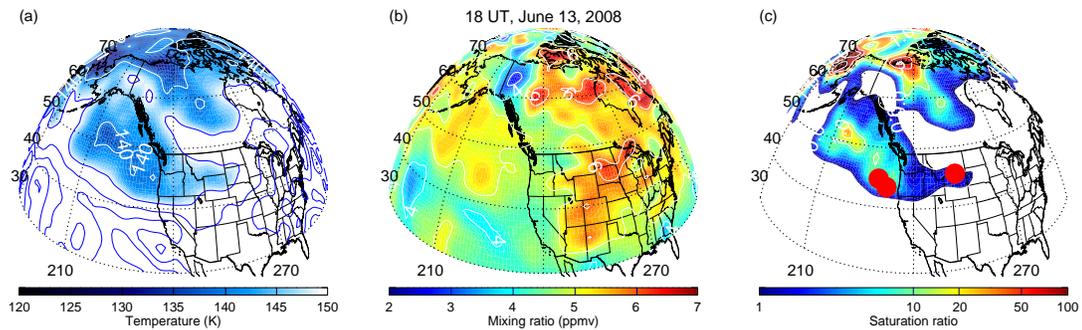


Fig. 18. Maps of analyzed NOGAPS-ALPHA (a) temperature (K), (b) water vapor mixing ratio (ppmv), and (c) saturation ratio  $S$  at 0.006 hPa on 13 June 2007 at 1800 UTC.

662 North American and Pacific regions on 13 June at 1800 UTC. These synoptic  
 663 maps show an outbreak of cold mid-latitude MLT air below 150 K over the  
 664 northern Pacific and northwestern US on this day. These low temperatures

665 extend to nearly 40°N and yield saturation ratios in excess of unity across the  
666 northwestern US and northeastern Pacific regions where the 3 cloud events in  
667 Figure 17 were detected. These comparisons show that the diagnostic estimate  
668 of saturation ratio  $S$  based on analyzed NOGAPS-ALPHA temperature and  
669 water vapor fields shows skill in identifying this mid-latitude burst of MCs ob-  
670 served by SHIMMER. The maps indicate that mid-latitude mesospheric clouds  
671 can exist here at this time due to the cold supersaturated local environment,  
672 without the need in this case for additional local subgrid-scale gravity wave  
673 temperature cooling (e.g., Rapp et al., 2002) or rapid equatorward advection  
674 of cloud particles formed at cooler regions to the north (e.g., Gerding et al.,  
675 2007).

## 676 *7.2 NLC Event of 19-20 June 2007*

677 In the late evening hours of 19 June 2007 and into the morning of 20 June,  
678 spectacular displays of noctilucent clouds (NLCs) were widely reported and  
679 photographed by many amateur observers in Washington state and Oregon in  
680 the western United States<sup>1</sup>, with bright NLCs photographed as far south as  
681 Bend, Oregon (44.0°N). The ground-based NLC observing network of Dalin et  
682 al. (2008) also reported very bright NLCs on this date at its western Canadian  
683 station.

684 Figure 19 plots the same sequence of temperature, water vapor and saturation  
685 ratios from the NOGAPS-ALPHA analysis of 20 June at 0600 UTC. The  
686 mesosphere over the northwestern US region is again cold, characterized by a  
687 tongue of cold air in Figure 19a that extends equatorward from polar regions

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<sup>1</sup> see, e.g., [http://www.spaceweather.com/nlcs/gallery2007\\_page4.htm](http://www.spaceweather.com/nlcs/gallery2007_page4.htm)

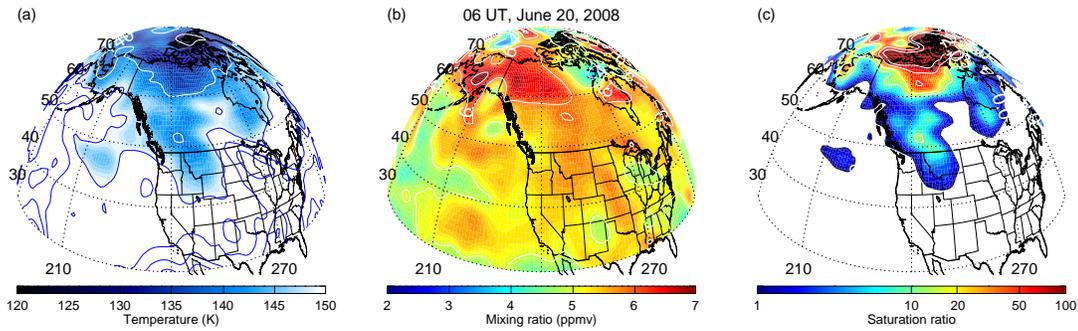


Fig. 19. Same presentation as Figure 18 but showing NOGAPS-ALPHA fields at 0.006 hPa on 20 June 2008 at 0600 UTC.

688 and within which moderate supersaturations occur in Figure 19c. Results on  
 689 1800 UTC (not shown) show this same feature has migrated westward to  
 690 lie over the Pacific in a similar band of latitudes. The diagnostic  $S$  values  
 691 from NOGAPS-ALPHA therefore appear consistent with these observed mid-  
 692 latitude NLCs.

### 693 7.3 NLC Event of 30 June 2007

694 NLCs were reported and photographed on 29-30 June 2007 from many regions  
 695 in Europe, including Sweden, Denmark, the United Kingdom, Hungary, and  
 696 even as far south as Portugal<sup>2</sup>.

697 The NOGAPS-ALPHA 0.006 hPa 0000 UTC analysis on 30 June 2007 in Fig-  
 698 ure 20 shows that an extensive pool of cold mesospheric air has moved over  
 699 northwestern Europe. Particularly interesting in this event are collars of en-  
 700 hanced water vapor mixing ratios in Figure 20b that wrap equatorward from  
 701 the eastern Atlantic and Eurasia into the United Kingdom and mid-latitude  
 702 European countries. These water vapor enhancements combine with low tem-

<sup>2</sup> see, e.g., [http://www.spaceweather.com/nlcs/gallery2007\\_page6.htm](http://www.spaceweather.com/nlcs/gallery2007_page6.htm)

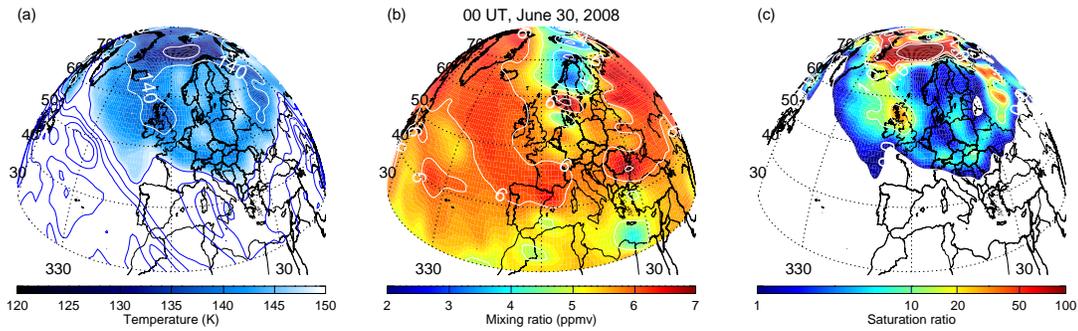


Fig. 20. Same presentation as Figure 18 but showing NOGAPS-ALPHA fields at 0.006 hPa on 30 June 2006 at 0000 UTC centered over western Europe.

703 peratures to yield large supersaturations in Figure 20c over, for example, the  
 704 United Kingdom and Hungary. Again, the widespread reports of NLC displays  
 705 over Europe at this time are borne out by enhanced mid-latitude saturation  
 706 ratios in Figure 20c over broad swaths of the continent, as derived from ana-  
 707 lyzed NOGAPS-ALPHA temperature and water vapor fields.

## 708 8 Summary and Outlook

709 We have presented first results of a specific high-altitude assimilation of MLS  
 710 and SABER temperatures and MLS water vapor and ozone mixing ratios us-  
 711 ing a new version of NOGAPS-ALPHA with a full DAS component extending  
 712 to MLT altitudes. These experiments were carefully designed, tuned and per-  
 713 formed during the northern summer MLT season from 15 May to 15 July,  
 714 2007, the first PMC season measured by instruments on the AIM satellite.

715 Assimilated temperatures from NOGAPS-ALPHA showed small biases with  
 716 respect to SABER, MLS, SOFIE and GEOS4 temperatures at most altitudes  
 717 and latitudes, including the polar summer MLT. Combining assimilated wa-

718 ter vapor and temperature fields into saturation ratios  $S$  yielded excellent  
719 agreement with seasonal variations in SOFIE-based PMC frequency, as well  
720 as several separate mid-latitude MC/NLC events observed by SHIMMER and  
721 ground-based observers at various times in June 2007.

722 Spectral signatures of the quasi 5-day (1,1) Rossby normal mode and solar  
723 migrating diurnal and semidiurnal tides were isolated from the analyses and  
724 investigated as a function of height, latitude, time and assimilated parameter.  
725 The 5-day wave temperature amplitudes at PMC altitudes of  $\sim 1-3$  K were  
726 similar to those inferred in previous studies. A clear 5-day wave signal in  
727 water vapor mixing ratios in the polar summer MLT is also revealed, but  
728 which does not correlate directly with the 5-day wave in temperature. Diurnal  
729 and semidiurnal tidal amplitudes appeared to be broadly realistic, despite the  
730 intrinsic limitations of the 6-hourly 3DVAR data insertion process.

731 These preliminary results underscore how emerging ground-to-MLT global  
732 analysis fields provided by this and other NWP systems (e.g., Polavarapu et  
733 al., 2005a) offer new opportunities for MLT science. Within the specific context  
734 of AIM and PMCs, these fields can be used to systematically investigate some  
735 of the major scientific issues, such as: hemispheric asymmetries in summer  
736 polar MLT temperatures and PMC amount (Siskind et al., 2003); planetary  
737 wave modulation of PMCs (von Savigny et al., 2007); relative roles of vertical  
738 and horizontal transport on mid-latitude PMCs (Gerding et al., 2007), and;  
739 possible global teleconnections between PMCs in one hemisphere and polar  
740 stratospheric meteorological conditions in the other hemisphere (Karlsson et  
741 al., 2007). In helping to answer these and other specific questions, these MLT  
742 fields can begin to address deeper questions about the fundamental charac-  
743 ter of the MLT itself. For example, how predictable is the MLT (Hoppel et

744 al., 2008)? Is MLT transport mainly local or nonlocal, and more specifically,  
745 is long-range horizontal transport in the MLT predictable or fundamentally  
746 chaotic (e.g., Holton and Schoeberl, 1988; Shepherd et al., 2000; Stevens et  
747 al., 2003)?

748 That having been said, current ground-to-MLT forecast-assimilation systems,  
749 including this one, are at the very early stages of their development and there is  
750 much left to do to improve them for future studies. In addition to assimilation  
751 of new data (e.g., SSMIS), new and improved model parameterizations and  
752 DAS algorithms are needed to handle the MLT region better (e.g., Polavarapu  
753 et al., 2005b; Sankey et al., 2007). Among the near-term foci for NOGAPS-  
754 ALPHA development are more completely tuned nonorographic GWD pa-  
755 rameterizations for all latitudes and seasons, more complete diabatic heating  
756 and cooling rate parameterizations for the MLT, parameterizations of water  
757 vapor and ozone chemistry in the MLT, NNMI algorithms that work well in  
758 the MLT, better MLT bias correction schemes, and higher-altitude radiative  
759 forward models  $\mathcal{H}$  for direct assimilation of MLT radiances into the system  
760 (Han et al., 2007). As the AIM mission continues through 2008, it will con-  
761 tinue to offer a superb testbed environment for developing and validating these  
762 capabilities, which in turn should hopefully feed back to add value to AIM  
763 measurements and enhance the science return of this mission.

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